

# Carrier Generation and Collection in CdS/CdSe-Sensitized SnO<sub>2</sub> Solar Cells Exhibiting Unprecedented Photocurrent Densities

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Semiconductor nanocrystals, especially quantum dots (QDs), have been attracting considerable attention due to their potential use as sensitizers in sensitized mesoscopic solar cells.<sup>1,2</sup> The ease with which the band gap can be tuned by varying particle size,<sup>3</sup> and other inherent absorption properties, makes QDs very promising candidates for efficient light harvesting materials in this type of solar cells.<sup>4,5</sup> In recent years, a lot of work on QD-sensitized solar cells has been reported where sensitization was achieved using CdS,<sup>6–11</sup> CdSe,<sup>12–14</sup> or PbS<sup>15–18</sup> and in some cases co-sensitization using more than one of these materials. In most reports, nanocrystalline TiO<sub>2</sub> is used as the photoanode.<sup>19–21</sup> The highest efficiencies of semiconductor-sensitized solar cells (SSCs) reported in the past few years are in the range of 4–5%.<sup>22–24</sup>

Nanostructured SnO<sub>2</sub> is known to have higher electron mobility than nanostructured TiO<sub>2</sub><sup>25,26</sup> and also a much more negative conduction band minimum (CBM).<sup>1,18,27</sup> Exploitation of the low CBM of SnO<sub>2</sub> should facilitate charge transfer from low band gap sensitizers such as near-infrared light absorbers, PbS,<sup>1,17</sup> PbSe,<sup>28</sup> and CuInSe<sub>2</sub>,<sup>29</sup> etc., where light harvesting could be significantly enhanced by utilizing long wavelength photons. Therefore, nanostructured SnO<sub>2</sub> is an interesting photoanode material, which could in principle facilitate efficient charge injection and charge collection. While the photovoltaic conversion efficiencies for cells employing SnO<sub>2</sub> are notoriously poor in dye-sensitized solar cells, it shows great promise in SSCs.<sup>30</sup>

It has been shown that the photoconversion efficiency of a QD-sensitized solar cell can be enhanced by co-sensitization of CdS and CdSe in a layered structure of TiO<sub>2</sub>/CdS/CdSe, where the band edges of the three

**ABSTRACT** CdS/CdSe-sensitized nanostructured SnO<sub>2</sub> solar cells exhibiting record short-circuit photocurrent densities have been fabricated. Under simulated AM 1.5, 100 mW cm<sup>-2</sup> illumination, photocurrents of up to 17.40 mA cm<sup>-2</sup> are obtained, some 32% higher than that achieved by otherwise identical semiconductor-sensitized solar cells (SSCs) employing nanostructured TiO<sub>2</sub>. An overall power conversion efficiency of 3.68% has been achieved for the SnO<sub>2</sub>-based SSCs, which compares very favorably to efficiencies obtained by the TiO<sub>2</sub>-based SSCs. The characteristics of these SSCs were studied in more detail by optical measurements, spectral incident photon-to-current efficiency (IPCE) measurements, and impedance spectroscopy (IS). The apparent conductivity of sensitized SnO<sub>2</sub> photoanodes is apparently too large to be measured by IS, yet for otherwise identical TiO<sub>2</sub> electrodes, clear electron transport features could be observed in impedance spectra, tacitly implying slower charge transport in TiO<sub>2</sub>. Despite this, electron diffusion length measurements suggest that charge collection losses are negligible in both kinds of cell. SnO<sub>2</sub>-based SSCs exhibit higher IPCEs compared with TiO<sub>2</sub>-based SSCs which, considering the similar light harvesting efficiencies and the long electron diffusion lengths implied by IS, is likely to be due to a superior charge separation yield. The resistance to charge recombination is also larger in SnO<sub>2</sub>-based SSCs at any given photovoltage, and open-circuit photovoltages under simulated AM 1.5, 100 mW cm<sup>-2</sup> illumination are only 26–56 mV lower than those obtained for TiO<sub>2</sub>-based SSCs, despite the conduction band minimum of SnO<sub>2</sub> being hundreds of millielectronvolts lower than that of TiO<sub>2</sub>.

**KEYWORDS:** semiconductor-sensitized solar cell · tin oxide · cadmium selenide · charge separation · charge collection · diffusion length

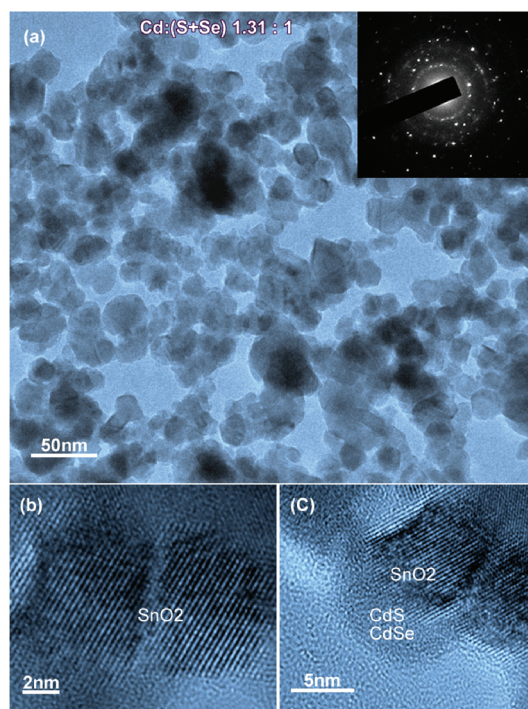
materials form a stepwise cascade, which is advantageous in the separation of excited electrons and holes across the interfacial region.<sup>22,31–34</sup> In our study, layers of semiconductors (CdS/CdSe/ZnS) were coated conformally onto SnO<sub>2</sub> by the successive ionic layer adsorption and reaction method (SILAR).<sup>22,31–36</sup> A passivation layer of TiO<sub>2</sub> on the SnO<sub>2</sub> surface is necessary before deposition of CdS in order to reduce the density of electron trap states at the surface of SnO<sub>2</sub>. Likewise, a barrier layer of ZnS, deposited by a post-treatment onto the core-shell CdS/CdSe, is required to inhibit the recombination of injected electrons with holes in the electrolyte, as investigated in a previous publication.<sup>30</sup>

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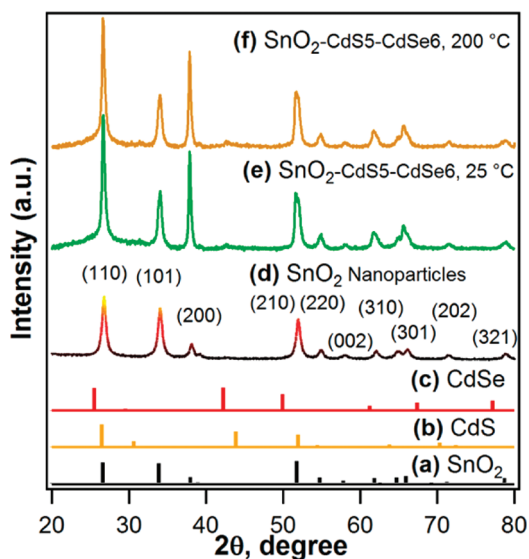


**Figure 1.** TEM image of a SnO<sub>2</sub>/CdS/CdSe electrode (a), HRTEM image of SnO<sub>2</sub> grains (b), and conformal coating of CdS/CdSe onto SnO<sub>2</sub> (c).

In this article, attempts to improve the photocurrent of SSCs by using SnO<sub>2</sub> as a photoanode instead of TiO<sub>2</sub> are described. Basic  $j$ - $V$ , incident photon-to-current conversion efficiency (IPCE), and optical measurements have been combined with impedance spectroscopy (IS) in order to scrutinize charge generation, transport, and recombination in these SSCs. Ultimately, a record photocurrent density of 17.40 mA cm<sup>-2</sup> was obtained under simulated AM 1.5, 100 mW cm<sup>-2</sup> illumination when an aqueous polysulfide electrolyte and a Cu<sub>2</sub>S cathode were used. The maximum IPCE was observed to be close to 80% over a wide wavelength range, indicating almost unity quantum efficiency, after correction for light absorption by the FTO substrate. On the other hand, for TiO<sub>2</sub>-based SSCs, a maximum IPCE of only 65% was obtained which, based upon results of IS, is thought to be due to inferior charge separation as opposed to charge collection. Very promising overall power conversion efficiencies of up to 3.68% were achieved, making nanostructured SnO<sub>2</sub> a potential rival to TiO<sub>2</sub> for use in SSCs.

## RESULTS AND DISCUSSION

**Morphology and Structural Examination.** Figure 1a shows a TEM image of CdS/CdSe-coated SnO<sub>2</sub> nanoparticles (5/5 SILAR cycles, *cf.* Experimental Section), with approximate sizes in the range of 20–40 nm. High-resolution transmission electron microscopy (HRTEM) suggests that these SnO<sub>2</sub> particles consist of crystal grains of ~12 nm (Figure 1b,c). Energy-dispersive X-ray spectroscopy (EDS) elemental analysis shows the



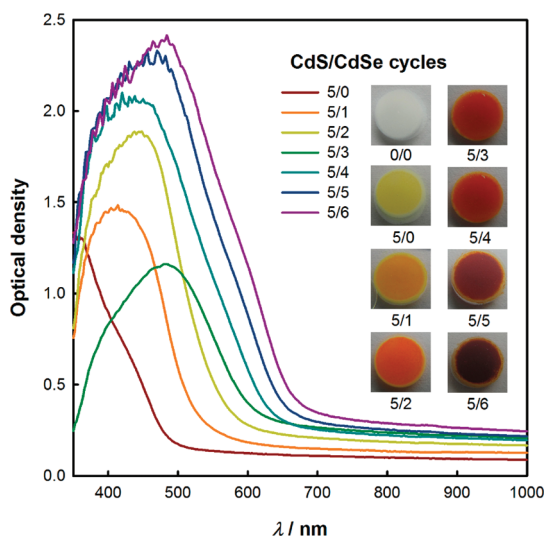
**Figure 2.** X-ray diffraction patterns of reference SnO<sub>2</sub> nanoparticles (a), CdS thin film (b), CdSe thin film (c) and measured diffraction patterns of SnO<sub>2</sub> nanoparticles (d), an untreated CdS/CdSe-coated SnO<sub>2</sub> sample (e), and after sintering in vacuum for 6 h at 200 °C (f).

existence of Cd, S, and Se. The atomic ratio Cd:(S+Se) was found to be 1.31:1, indicating that the layer as a whole was Cd-rich and broadly consistent with previous reports where Cd-rich CdSe was deposited by SILAR.<sup>14</sup> A conformal coating of CdS/CdSe on the SnO<sub>2</sub> was revealed by HRTEM (Figure 1c), and the lack of lattice fringes suggests the as-grown sensitizer is amorphous.

The crystal structure of samples was characterized by powder X-ray diffraction (XRD). As shown in Figure 2d, the XRD pattern revealed that the SnO<sub>2</sub> nanoparticles are polycrystalline, consistent with the HRTEM image (Figure 1b,c), which shows lattice fringes with  $d$  spacing of 0.335 nm, corresponding to the lattice plane (110) of tetragonal rutile SnO<sub>2</sub> (JCPDS No. 41-1445, Figure 2a). The calculated grain size of ~12 nm from the Scherrer equation using the first and second peaks (Figure 2d) was consistent with the grain size obtained from analysis of HRTEM images. Additional features due to CdS and CdSe were not observed in the X-ray diffraction analysis, except for a noisier profile (Figure 2e,f), arising because of the amorphous nature of the conformal coating of CdS/CdSe.

### Optical Properties of CdS/CdSe-Sensitized SnO<sub>2</sub> Electrodes.

Nanocrystalline CdS and CdSe are widely used semiconductor sensitizers for nanocrystalline TiO<sub>2</sub> solar cells because their composite absorption spectrum covers most of the visible region.<sup>22,31,37</sup> The SILAR method offers the easiest way to deposit semiconductor absorber layers onto nanocrystalline electrodes, where mostly the coating is conformal. Co-sensitization by core-shell quantum dots yields important absorption properties and charge separation characteristics, which depend on the relative band positions of the constituent sensitizer materials.<sup>30</sup>



**Figure 3.** UV–vis optical density spectra of CdS and CdS/CdSe-coated SnO<sub>2</sub> electrodes alongside typical photographs of electrodes with or without five cycles of CdS and different numbers of CdSe cycles.

Figure 3 shows optical density (OD) spectra of a 6.4 μm thick SnO<sub>2</sub> electrode as it is coated with CdS/CdSe by successive SILAR cycles. The incorporated amount of CdS or CdSe into the SnO<sub>2</sub> films increases with increasing numbers of SILAR cycles. As shown in Figure 3, the apparent absorption onset of five-cycle CdS/SnO<sub>2</sub> electrodes occurs at ~500 nm, representing its bulk nature and optical band gap of 2.50 eV.<sup>38</sup> Successive color changes from yellow ochre (5/1) to dark red (5/6) were observed with the increase of CdSe deposition cycles (inset of Figure 3), indicating strong absorption of most of the visible wavelengths.

For all spectra, an apparent absorbance extending up to 1000 nm was observed which systematically increases with the number of SILAR deposition cycles. We suspect this is due to uncorrected light scattering losses which increase with increasing CdS/CdSe coverage; rigorous corrections for light scattering using diffuse transmittance and reflectance measurements have not yet been attempted. Despite the ambiguities caused by scattering effects, it still seems very reasonable to assume that the light harvesting efficiency of the sensitized electrodes is almost unity in the wavelength range of 400–500 nm, where an OD of >2 is found for *ca.* 6 μm thick electrodes (note that the electrodes used in complete SSCs were 9 μm thick). The same conclusion can be drawn for TiO<sub>2</sub> electrodes sensitized with 5/5 CdS/CdSe SILAR cycles, which have similar OD spectra (Supporting Information).

A deviation in the magnitude of OD was also observed for the specific case of a 5/3 combination of CdS/CdSe, where the absorption onset followed the trend of the other films, but the peak absorption was less than for 5/2 and 5/4 CdS/CdSe combinations. This surprising result was reproducible; adding an

additional CdSe cycle to a 5/2 electrode resulted in a decrease in peak OD together with the expected extension of the absorption onset to longer wavelengths. Since the detailed composition and morphology of these CdS/CdSe layers is unknown, we do not wish to speculate about the origin of this peculiar phenomenon, which is the subject of ongoing investigation in our laboratory.

These co-sensitized electrodes have wide wavelength absorption characteristics, emphasizing that the main sensitization comes from CdSe, due to its wider spectral response than CdS. Likewise, due to size-induced electronic quantization phenomena, the absorption onset of CdS/CdSe-sensitized electrodes moves to longer wavelengths with the increase of CdSe deposition cycles, indicating the increment of the thickness of the CdSe layer. The performance of SSCs incorporating CdS/CdSe-sensitized electrodes has been studied, and the optimum layer combination is found to be 5/5 CdS/CdSe. It should be mentioned that the samples of SnO<sub>2</sub>/CdS/CdSe used for OD measurements were prepared without TiCl<sub>4</sub> pretreatment of the SnO<sub>2</sub> electrodes or the ZnS post-treatment. It has been reported that the ZnS layer decreases quantum confinement and red shifts the spectral absorption, but hardly participates in harvesting of incident light itself, rather it plays an important role in improving charge separation and collection.<sup>31</sup>

**Photocurrent–Voltage and Energy Conversion Characteristics.** To study the photovoltaic performance of CdS/CdSe-sensitized mesoscopic SnO<sub>2</sub> electrodes, sandwich-type thin layer cells were fabricated with FTO/SnO<sub>2</sub>/CdS/CdSe as the photoanode, platinumized FTO and Cu<sub>2</sub>S on brass as the cathode, and a polysulfide electrolyte as the hole transporter. To prevent recombination at the SnO<sub>2</sub>/CdS/CdSe/electrolyte interface, SnO<sub>2</sub> electrodes were treated with TiCl<sub>4</sub> aqueous solution prior to CdS/CdSe deposition, and a thin passivation layer of ZnS was deposited onto the sensitized electrode by two SILAR deposition cycles.

Figure 4a shows the IPCE spectra of CdS/CdSe-sensitized SnO<sub>2</sub> and TiO<sub>2</sub> solar cells, which exhibit a strong photoresponse over the entire visible light range, with photon wavelengths even extending to ~850 nm for the SnO<sub>2</sub> SSCs. It has been shown that, for the co-sensitized electrode, the higher IPCE in the short-wavelength region is due to the CdS, whereas CdSe harvests light in the long wavelength region; together a better overall IPCE response was observed.<sup>39</sup> IPCEs of ~80% between 400 and 600 nm were achieved with 5/5 layers of CdS/CdSe on SnO<sub>2</sub> and 2 layers of ZnS. Since the IPCE is found to be within a few percent of the transmission of the FTO substrate in this wavelength range, it must be concluded that light harvesting, charge separation, and charge collection are all close to unity when SnO<sub>2</sub> electrodes are used.



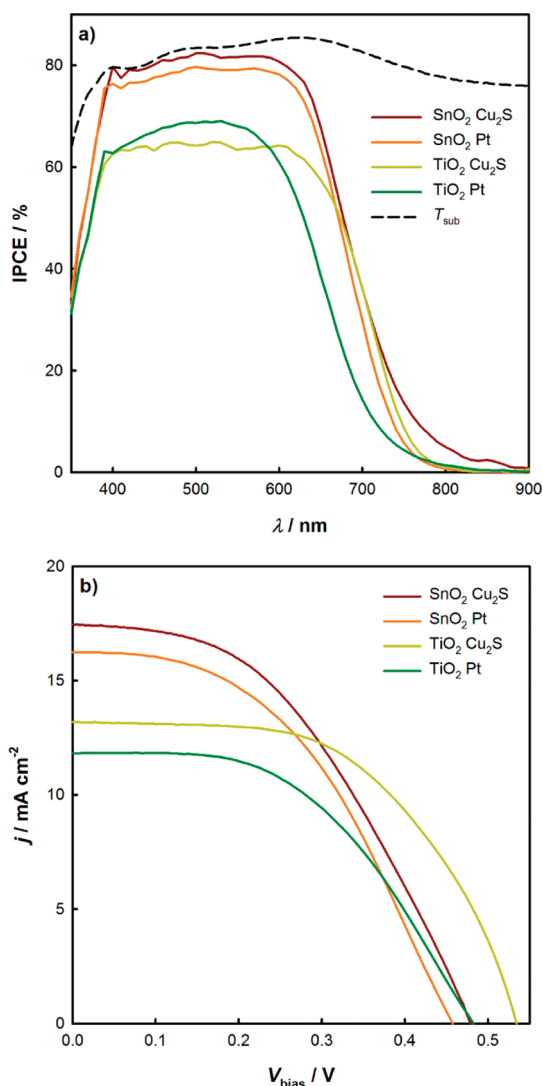


Figure 4. IPCE spectra (a) and  $j$ - $V$  characteristics (b) of CdS/CdSe-sensitized SnO<sub>2</sub> and TiO<sub>2</sub> solar cells.

**TABLE 1. Characteristics of CdS/CdSe-Sensitized SnO<sub>2</sub> and TiO<sub>2</sub> Solar Cells with Platinized FTO and Cu<sub>2</sub>S Cathodes under Simulated AM 1.5, 100 mW cm<sup>-2</sup> Illumination**

photoanode	CdS/CdSe	cathode	$V_{oc}$ (V)	$j_{sc}$ (mA cm <sup>-2</sup> )	ff (%)	$\eta$ (%)
SnO <sub>2</sub>	5/5	Pt	0.458	16.20	45.85	3.43
	5/5	Cu <sub>2</sub> S	0.477	17.40	44.40	3.68
TiO <sub>2</sub>	5/5	Pt	0.484	11.82	49.39	2.85
	5/5	Cu <sub>2</sub> S	0.533	13.17	55.28	3.88

Figure 4b shows the  $j$ - $V$  characteristics of sensitized SnO<sub>2</sub> and TiO<sub>2</sub> SSCs. Photovoltaic parameters of the SSCs are given in Table 1. As expected from the IPCE measurements, SnO<sub>2</sub> electrodes exhibit superior short-circuit photocurrents as compared to TiO<sub>2</sub> electrodes. The short-circuit photocurrents ( $j_{sc}$ ) predicted by convolution of IPCE spectra with the AM 1.5 spectrum followed by integration over all wavelengths underestimate the measured 1 Sun  $j_{sc}$  values by *ca.* 10%, implying differences in charge separation and charge collection

between the IPCE measurements (which were made using low intensity monochromatic illumination) and the 1 Sun  $j$ - $V$  measurements are minimal.

The best SnO<sub>2</sub> device was obtained with 5/5 CdS/CdSe as well as TiCl<sub>4</sub> and ZnS pre- and post-treatments, respectively, yielding  $j_{sc} = 17.4 \text{ mA cm}^{-2}$ ,  $V_{oc} = 0.477 \text{ V}$ , fill factor = 44.4%, and  $\eta = 3.68\%$  under simulated AM 1.5, 100 mW cm<sup>-2</sup> illumination. One plausible reason for the higher photocurrent of the SnO<sub>2</sub> SSC compared with the TiO<sub>2</sub> SSC is better charge transport dynamics in SnO<sub>2</sub> arising from the higher electron mobility, which will shortly be discussed in the light of results from electrochemical impedance spectroscopy experiments. It is also known that SnO<sub>2</sub> has a much lower conduction band minimum (CBM) than TiO<sub>2</sub> in the same media. The large driving force for charge injection originating from the low CBM of SnO<sub>2</sub> may help efficient charge injection from excited CdS/CdSe to the conduction band of SnO<sub>2</sub>. Leventis *et al.* recently reported comparative studies of charge separation between PbS-sensitized mesoscopic TiO<sub>2</sub> and SnO<sub>2</sub>.<sup>17</sup> Transient optical measurements showed much faster charge injection for the PbS-sensitized SnO<sub>2</sub> electrode.<sup>17</sup> Guijarro *et al.* have shown that the injection of electrons from CdSe into TiO<sub>2</sub> decreases significantly with increasing numbers of SILAR cycles.<sup>14</sup> This could be eliminated by using a low CBM oxide such as SnO<sub>2</sub> for an electron acceptor material. Pijpers *et al.* have shown efficient electron injection from PbSe into SnO<sub>2</sub>; however, there was no injection into TiO<sub>2</sub>, which was attributed to the different energy band alignments when using these oxides.<sup>40</sup>

As found previously,<sup>30</sup> a TiCl<sub>4</sub> treatment effectively passivates the surface of SnO<sub>2</sub> and significantly improves the photovoltaic performance of CdSe-sensitized SnO<sub>2</sub> solar cells. Prasittichai *et al.* have found that the photovoltaic performance of dye-sensitized SnO<sub>2</sub> solar cells increased 5-fold when the surface was passivated with an ultrathin layer of Al<sub>2</sub>O<sub>3</sub>, stemming from inhibited recombination.<sup>41</sup> In addition, SILAR deposition forms a conformal coating of CdS and CdSe on SnO<sub>2</sub>, which acts as a barrier layer (in a similar way to Al<sub>2</sub>O<sub>3</sub>), preventing injected electrons in SnO<sub>2</sub> from recombining with holes in the electrolyte. Moreover, it is believed that surface modification by ZnS further inhibits the recombination of excited electrons at the electrode/electrolyte interface and increases the injection from the sensitizer into TiO<sub>2</sub>.<sup>42-44</sup>

In this work, optimized TiO<sub>2</sub> SSCs yielded  $j_{sc} = 13.17 \text{ mA cm}^{-2}$ ,  $V_{oc} = 0.533 \text{ V}$ , fill factor = 55.28%, and  $\eta = 3.88\%$  when Cu<sub>2</sub>S cathodes were used. The better overall performance of the TiO<sub>2</sub> cell with the same configuration arises from the higher fill factor and slightly higher  $V_{oc}$ , which compensate for the lower photocurrent. When platinized cathodes were used in place of Cu<sub>2</sub>S cathodes, performance with both

oxides was found to be lower due to inferior fill factor and  $j_{sc}$ . While the photocurrent of CdS/CdSe-sensitized SnO<sub>2</sub> SSCs is significantly higher than the value reported for CdS/CdSe-sensitized TiO<sub>2</sub> cells,<sup>19,22,31</sup> the overall efficiency still remains lower. It is not surprising that the  $V_{oc}$  of the SnO<sub>2</sub> cell would be lower than the TiO<sub>2</sub> SSCs because of the low CBM and unchanged redox level of the polysulfide electrolyte, and in fact, a much larger loss in  $V_{oc}$  might be expected.

The effect of the SnO<sub>2</sub> mesoporous film thickness on the cell performance was investigated by varying the film thickness from 5.5 to 10.5  $\mu\text{m}$ , with the same CdS/CdSe sensitizer loading. Photocurrents normally increased with film thickness, but  $V_{oc}$  and fill factor dropped, eventually leading to a decrease in efficiency for thicknesses  $>9 \mu\text{m}$ . For the TiO<sub>2</sub> SSCs, an optimum thickness of  $\sim 9 \mu\text{m}$  was also found.

**Investigation of Charge Transport and Recombination Processes Using Impedance Spectroscopy.** Impedance spectroscopy (IS) was used to compare the charge transport and charge transfer processes occurring in TiO<sub>2</sub>- and SnO<sub>2</sub>-based SSCs. Figure 5 shows representative IS spectra for TiO<sub>2</sub>- (Figure 5a) and SnO<sub>2</sub>-based SSCs (Figure 5b). For the case of the TiO<sub>2</sub>-based SSCs, excellent fits to most spectra (dashed red lines in Figure 5a) can be obtained using an equivalent circuit (Supporting Information) based upon the now ubiquitous diffusion-reaction model for a porous electrode, as described by Bisquert.<sup>45</sup>

In spectra recorded for TiO<sub>2</sub>-based SSCs, for photovoltages in the range of 0.1–0.3 V, a clear Warburg-like feature can be observed, which is consistent with previous reports and can be assigned to electron diffusion in the TiO<sub>2</sub> layer.<sup>22</sup> The associated electron transport resistance ( $R_t$ ) decreases as open-circuit photovoltage ( $V_{oc}$ ) becomes increasingly more negative, corresponding to the electron quasi-Fermi level rising toward the CBM and the free carrier concentration in the semiconductor increasing. The slope on a semilogarithmic plot of  $R_t$  versus  $V_{oc}$  (dashed black line in Figure 5c) is  $16.6 \text{ V}^{-1}$ , which is in good agreement with the theoretical value of  $17 \text{ V}^{-1}$  at 295 K, lending support to the assignment of the Warburg-like impedance feature to the electron transport process.

On the contrary, for the SnO<sub>2</sub>-based SSCs, no high-frequency Warburg-like feature is observed under any conditions (*cf.* Figure 5b), and excellent fits to the data can be obtained with a simplified equivalent circuit consisting of the series connection of a resistor and two parallel RC circuits. One RC circuit represents the SnO<sub>2</sub> layer (with negligible resistance to electron transport), and the other represents the cathode–electrolyte interface. This finding implies that the electron transport resistance for these cells is always much lower than the cathode impedance; if

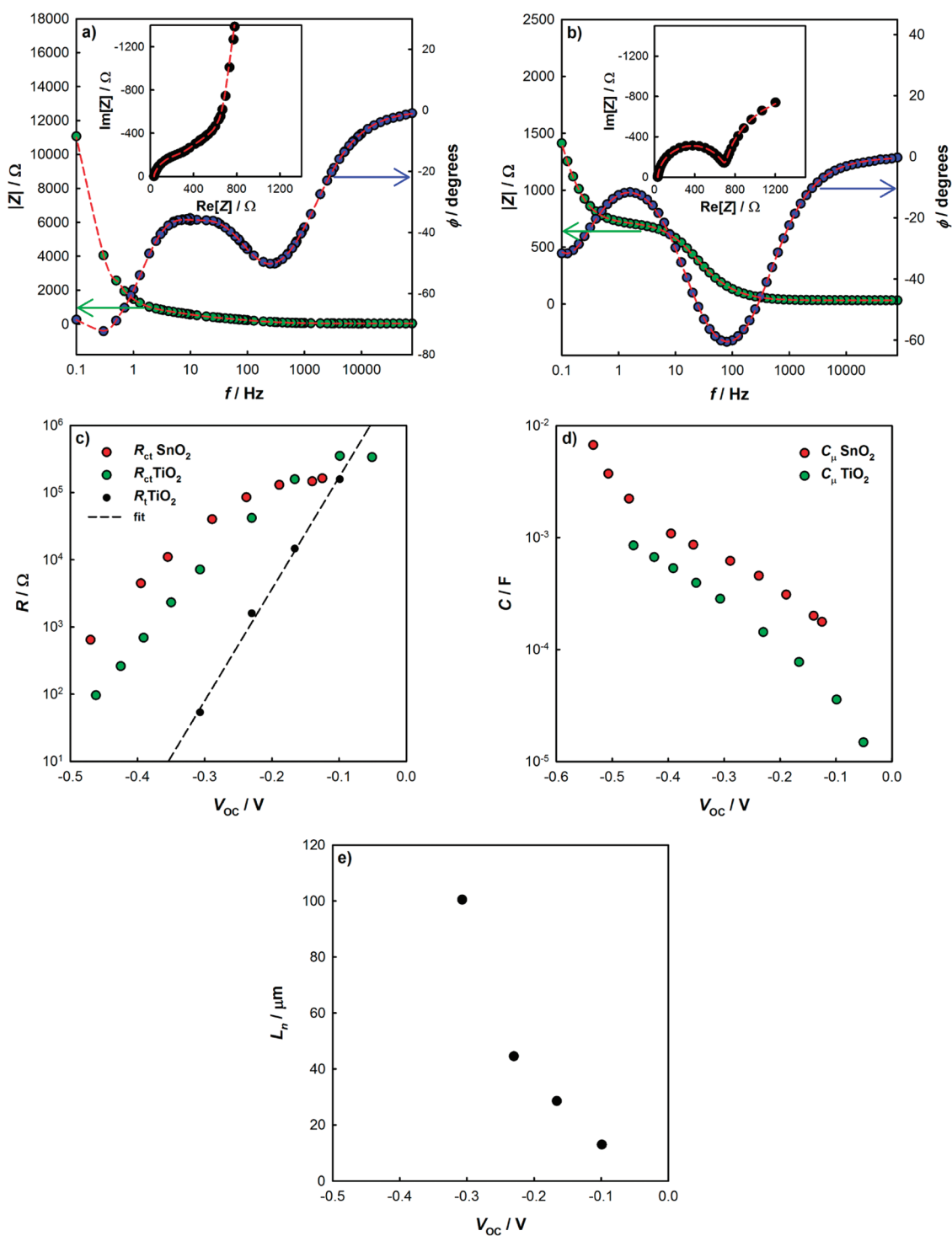
this were not the case, it ought to be apparent in the impedance spectrum. Consequently, the electron diffusion length in these cells is expected to be exceedingly long compared with the SnO<sub>2</sub> layer thickness.

A crude estimate of the minimum electron diffusion length at the cell voltage of the maximum power point ( $V_{mpp}$ ) for the SnO<sub>2</sub>-based SSC can be made by taking the SnO<sub>2</sub> interfacial charge transfer resistance ( $R_{ct}$ ) at a photovoltage corresponding to  $V_{mpp}$  and assuming the electron transport resistance ( $R_t$ ) is equal to the cathode charge transfer resistance (probably a gross overestimation). This approach yields  $L_n$  of the order of 10 times the SnO<sub>2</sub> layer thickness, implying a collection efficiency of practically unity, regardless of optical considerations. This estimate is entirely consistent with the IPCE spectra for SnO<sub>2</sub>-based SSCs employing Cu<sub>2</sub>S cathodes, where IPCE values approaching the substrate transmission are found (*cf.* Figure 4a).

Attempts were made to more precisely estimate  $L_n$  at various photovoltages by using a combination of intensity-modulated photovoltage spectroscopy (IMVS) and intensity-modulated photocurrent spectroscopy (IMPS). However, ultimately, this was not possible because under all conditions studied (open-circuit, various light intensities) the IMPS response was found to be dominated by the RC time constant of the cells and deconvolution of the transport time constant with any reasonable degree of accuracy was impossible.

Attempts were also made to observe electron transport features in the IS spectra by using a Cu<sub>2</sub>S electrode to reduce the impedance of the cathode. However, due to the instability of the Cu<sub>2</sub>S cathodes used in this work (which were fabricated by reaction of a brass substrate with the electrolyte), interpretation of impedance spectra were complicated by non-negligible drift of the system on the time scale of the measurements (which typically took 1–2 h), thus further analysis will not be presented here. However, it is noteworthy that, in these experiments, features in the impedance spectra (*e.g.*, slope of  $\sim 1$  in Nyquist plots) which were clearly attributable to an electron transport resistance were not observed.

For the TiO<sub>2</sub>-based SSCs, long electron diffusion lengths are also found for  $V_{oc} > 0.2 \text{ V}$ , where  $L_n$  is greater than 3 times the TiO<sub>2</sub> layer thickness, sufficient to ensure almost 100% efficient charge collection, regardless of illumination direction and absorption coefficient.<sup>46</sup> Below  $V_{oc} = 0.2 \text{ V}$ ,  $L_n$  decreases into a regime where charge collection losses could begin to become a photocurrent loss mechanism (the extent of which will depend upon the precise absorption coefficient of the sensitized metal oxide layer and on the illumination direction). However, this exceptionally low quasi-Fermi level position (obtained



**Figure 5.** Bode and Nyquist plots showing typical IS results and best fits using the diffusion-reaction model (dashed red lines) for TiO<sub>2</sub>- (a;  $V_{oc} = -0.23$  V, incident photon flux  $ca. 3 \times 10^{14} \text{ cm}^{-2} \text{ s}^{-1}$ ,  $\lambda = 627$  nm) and SnO<sub>2</sub>- (b;  $V_{oc} = -0.19$  V, incident photon flux  $ca. 3 \times 10^{14} \text{ cm}^{-2} \text{ s}^{-1}$ ,  $\lambda = 627$  nm) based SSCs; charge transfer resistance and transport resistance (c), and chemical capacitance (d) parameters obtained from fitting; dependence of electron diffusion length derived from IS fitting results on open-circuit photovoltage for a TiO<sub>2</sub>-based SSC (e).

under very low incident light intensity) is unlikely to be encountered in the device under normal 1 Sun operating conditions, where the average quasi-Fermi level in the bulk of the metal oxide layer for cell voltages below  $V_{mpp}$  (e.g., short-circuit) is expected to be only 0.15–0.25 V lower than  $V_{oc}$  at the same light

intensity.<sup>47–49</sup> It is noteworthy that the electron mobility in single crystal rutile SnO<sub>2</sub><sup>25</sup> is  $ca.$  10 times higher than in single-crystal anatase TiO<sub>2</sub>,<sup>26</sup> and it is possible that this is a factor contributing to the conspicuous absence of a prominent Warburg feature in the IS spectra for SnO<sub>2</sub>-based SSCs, despite

it being observed for otherwise identical TiO<sub>2</sub>-based SSCs.

Peak IPCE and  $j_{sc}$  values for SnO<sub>2</sub>-based SSCs are around 30–40% higher than for TiO<sub>2</sub>-based SSCs. Upon the basis of the diffusion length estimations for the TiO<sub>2</sub>-based device, it appears unlikely that this difference can be explained by charge collection losses in the TiO<sub>2</sub>-based SSCs. This is because, despite the apparently superior transport properties of SnO<sub>2</sub> compared with TiO<sub>2</sub>, transport is still *fast enough* to avoid collection losses in the TiO<sub>2</sub>-based SSCs. Assuming that differences in light harvesting efficiency between SnO<sub>2</sub> and TiO<sub>2</sub> electrodes are not significant over most of the absorbed wavelength range (this assumption is supported by transmission measurements for  $\lambda = 400$  to at least  $\lambda = 500$  nm; Supporting Information), it is likely that the cause for higher IPCE and photocurrent is that the overall charge separation yield in SnO<sub>2</sub>-based SSCs is higher than in TiO<sub>2</sub>-based SSCs. Since identical electrolytes and sensitizer deposition procedures were used for both types of SSC, it must be concluded that either charge injection from the CdS/CdSe to the metal oxide is slower in the TiO<sub>2</sub> system compared to the SnO<sub>2</sub> system, or that the reverse process, geminate recombination, is faster in the TiO<sub>2</sub> system.

The CBM of SnO<sub>2</sub> is thought to lie 0.3–0.5 eV below that of anatase TiO<sub>2</sub>, and therefore, on purely energetic grounds, a higher charge injection efficiency might be expected. However, one must also consider that use of Marcus–Gerischer electron transfer theory is probably more appropriate for predicting the probability of electron transfer from the semiconductor sensitizer to the metal oxide. In this case, the density of acceptor states at the relevant energy (the energy of the excited states of CdS/CdSe) in the metal oxide is more important than the free energy difference between initial and final states in the charge transfer process.

The conventional interpretation of the porous electrode capacitance obtained from IS data using the diffusion-reaction model is that it corresponds to a chemical capacitance, that is, a measure of the effective density of states around the quasi-Fermi level, which, in turn, is related to  $V_{oc}$  by  $V_{oc} = nE_F - E_{F,redox}$ . Since identical electrolytes were used in both types of SSC, comparisons of  $C_{\mu}$  at matched  $V_{oc}$  ought to reveal differences between the densities of states in the two oxides, on an energy reference scale relevant to the operation of these cells. Figure 5d clearly shows that, at any given  $V_{oc}$ ,  $C_{\mu}$  for the SnO<sub>2</sub>-based device is at least 2 times larger than for the TiO<sub>2</sub>-based device, and at lower voltages is almost 4 times larger. Thus, if it is permissible to extrapolate these results to the energies relevant to charge injection, the higher IPCE values for the SnO<sub>2</sub>-based SSCs

compared with the TiO<sub>2</sub> counterparts may be due to enhanced electron injection, as a result of the higher density of electron acceptor states at the appropriate energy. The shift along the voltage axis between the two plots in Figure 5d might also be viewed as evidence that the SnO<sub>2</sub> layer has a lower conduction band minimum than TiO<sub>2</sub>, as expected. However, since the density of states functions for the two materials are not known and are unlikely to be the same, quantifying the difference in band edge energies is not possible without more information.

It is also interesting to note that the charge transfer resistance ( $R_{ct}$ ) for the SnO<sub>2</sub>-based SSC is up to 6 times larger than that for the TiO<sub>2</sub>-based SSC. At very low photovoltages, similar  $R_{ct}$  values are found, but data obtained at these photovoltages are not likely to be relevant to normal operation of the cells. The larger  $R_{ct}$  is a factor that contributes toward the surprisingly high photovoltages achieved by the SnO<sub>2</sub>-based SSCs, which are only 26–56 mV lower than for TiO<sub>2</sub>-based SSCs, despite the far more positive CBM of SnO<sub>2</sub>. However, it must also be borne in mind that the effective density of states at the CBM ( $N_C$ ), together with the energy of the CBM, determines the free electron concentration required to produce a particular quasi-Fermi level position and  $V_{oc}$ . The effective electron mass in SnO<sub>2</sub> is reported to be *ca.* 35 times smaller than estimates for nanocrystalline TiO<sub>2</sub>,<sup>50,51</sup> thus significant differences in  $N_C$  between the two materials probably exist, and the relatively small difference in  $V_{oc}$  between the cells cannot be solely attributed to slower charge recombination in the SnO<sub>2</sub>-based SSCs.

## CONCLUSIONS

CdS/CdSe-sensitized SnO<sub>2</sub> solar cells exhibiting record photocurrent densities of up to 17.40 mA cm<sup>-2</sup> under simulated AM 1.5, 100 mW cm<sup>-2</sup> illumination have been fabricated. Consistent with the high photocurrent, an IPCE of close to 80% over almost the entire visible spectral region is observed, also implying highly efficient charge separation and collection. Open-circuit photovoltage is similar to that obtained by TiO<sub>2</sub>-based SSCs, despite the much more positive CBM of SnO<sub>2</sub>. Impressive overall power conversion efficiencies of up to 3.68% were achieved, making nanostructured SnO<sub>2</sub> a rival to TiO<sub>2</sub> for use in SSCs.

A detailed comparison between SnO<sub>2</sub> and TiO<sub>2</sub> photoanodes reveals SnO<sub>2</sub> can be considered a superior material to TiO<sub>2</sub> in many respects, owing to higher recombination resistances, faster charge transport, and more efficient charge separation. However, the efficiency of SSCs employing SnO<sub>2</sub> may be limited by the relatively low conduction band energy compared with TiO<sub>2</sub>. To overcome this problem, we anticipate



that control of surface dipoles at the various interfaces (e.g., SnO<sub>2</sub>/CdS, ZnS/electrolyte, etc.) will be important in engineering band edge positions with respect to the electrolyte redox level and consequently improving  $V_{oc}$ . Another strategy to improve conversion efficiency might involve replacement of

the polysulfide electrolyte with one possessing a more positive redox potential; however, this approach will almost certainly lead to changes in sensitizer regeneration rate and charge recombination rate, which may or may not be beneficial to device performance.

## EXPERIMENTAL SECTION

**Preparation of Mesoscopic SnO<sub>2</sub> Electrodes.** Tin(IV) oxide (SnO<sub>2</sub>) nanoparticles (Alfa Aesar, Nanoarc) mixed with a small amount of ethanol were ground by pestle and mortar to break up agglomerates. Ethyl cellulose and  $\alpha$ -terpineol were then mixed with the ground SnO<sub>2</sub> powder following a previously reported procedure for making TiO<sub>2</sub> paste.<sup>52</sup> SnO<sub>2</sub> electrodes were prepared by screen-printing the above paste onto FTO glass (TEC, 15  $\Omega/\square$ ) several times in order to get an appropriate thickness. The printed films were then sintered in air by heating gradually to 325 °C and holding for 5 min, then 375 °C for 5 min, at 450 °C for 15 min, and finally at 500 °C for 15 min. The resulting mesoporous SnO<sub>2</sub> electrodes were semi-transparent. The thickness of the electrodes was determined by an Alpha-Step IQ surface profiler to be  $\sim 9 \mu\text{m}$ . After preparation, SnO<sub>2</sub> electrodes were immersed into 40 mM aqueous TiCl<sub>4</sub> solution at 70 °C for 40 min and then washed with water and ethanol followed by drying in an electric oven at 70 °C.

**Preparation of CdS/CdSe-Sensitized Mesoscopic SnO<sub>2</sub> Electrodes.** After masking the bare FTO with Kapton tape, leaving the mesoscopic SnO<sub>2</sub> electrodes (6 mm diameter) uncovered, they were sensitized with CdS and CdSe using previously reported SILAR methods.<sup>19,22,30</sup> In a typical procedure, the SnO<sub>2</sub> electrodes were immersed in a solution containing 0.02 M cadmium nitrate tetrahydrate (Cd(NO<sub>3</sub>)<sub>2</sub>·4H<sub>2</sub>O, Fluka, >99.0%) in methanol for 1 min, to allow Cd<sup>2+</sup> to adsorb onto the SnO<sub>2</sub>, and then rinsed with methanol for 1 min to remove the excess Cd<sup>2+</sup>. Electrodes were then dried in a gentle stream of N<sub>2</sub> for 1 min. The dried electrodes were then dipped into a solution containing 0.02 M sodium sulfide nonahydrate (Na<sub>2</sub>S·9H<sub>2</sub>O, Sigma Aldrich) in a mixture of methanol and deionized water (1:1, v/v) for 1 min, where the preadsorbed Cd<sup>2+</sup> reacts with S<sup>2-</sup> to form the desired CdS. Electrodes were then rinsed in methanol for 1 min and dried again with N<sub>2</sub>. This procedure was repeated five times in order to get a suitable CdS loading on the SnO<sub>2</sub> electrode. CdSe was also deposited onto the CdS-coated SnO<sub>2</sub> electrodes by the SILAR method, where they were dipped into a solution containing 0.03 M cadmium nitrate tetrahydrate (Cd(NO<sub>3</sub>)<sub>2</sub>·4H<sub>2</sub>O, Fluka, >99.0%) in ethanol for 30 s and rinsed with ethanol for 2 min, then dried for 2 min in an argon atmosphere. Subsequently, the dried electrodes were dipped in to a solution containing 0.03 M Se<sup>2-</sup> for 30 s. The Se<sup>2-</sup> solution was prepared by mixing selenium dioxide (SeO<sub>2</sub>, Sigma-Aldrich, 99.9%) and sodium borohydride (NaBH<sub>4</sub>, Sigma Aldrich) in ethanol following the work reported by Lee *et al.*<sup>19</sup> One deposition cycle was completed by further rinsing in ethanol for 2 min and drying again in an argon atmosphere for 2 min. This procedure was repeated several times to get suitable CdSe loading on CdS-coated SnO<sub>2</sub> electrodes. A ZnS passivation layer was deposited on a CdS/CdSe-coated SnO<sub>2</sub> electrode by SILAR using 0.1 M aqueous zinc acetate dihydrate (Zn(CH<sub>3</sub>COO)<sub>2</sub>·2H<sub>2</sub>O, Sigma Aldrich, >99.5%) and 0.1 M aqueous Na<sub>2</sub>S·9H<sub>2</sub>O, dipping for 1 min in each solution, with 1 min of rinsing in deionized water between each immersion in a precursor solution. All the CdS/CdSe-sensitized SnO<sub>2</sub> electrodes used in this study were coated with two ZnS passivation layers. For comparison, TiO<sub>2</sub> electrodes were prepared by screen-printing TiO<sub>2</sub> paste (Degussa P25 powder and ethyl cellulose in  $\alpha$ -terpineol) on bare FTO to produce  $\sim 9 \mu\text{m}$  thick layers. The rest of the fabrication procedure was identical to that used for CdS/CdSe-sensitized SnO<sub>2</sub> electrodes.

**Material Characterization.** Powder X-ray diffraction (XRD) measurements of SnO<sub>2</sub> nanoparticles and SnO<sub>2</sub>/CdS/CdSe electrodes were carried out with a Bruker D8 using Cu K $\alpha$ 1 radiation ( $\lambda = 0.154059 \text{ nm}$ ). The detailed morphology of the SnO<sub>2</sub> particles were elucidated from transmission electron microscopy on JEOL JEM 2010F, which also facilitated energy-dispersive X-ray analysis (EDX). Samples for TEM investigations were prepared by detaching the CdS/CdSe-sensitized SnO<sub>2</sub> nanocrystals mechanically from the FTO glass substrate and dispersing in ethanol, followed by transferring one drop of the suspension onto a carbon-coated copper grid. Optical density/transmission spectra were collected using a Shimadzu UV-vis-NIR spectrophotometer (SolidSpec-3700) at room temperature.

**Fabrication and Characterization of CdS/CdSe-Sensitized SnO<sub>2</sub> Solar Cells.** Platinized counter electrodes were fabricated on FTO with small holes drilled into one corner. After cleaning, a thin layer of Pt was deposited onto the FTO by thermal decomposition of hexachloroplatinic acid. Brass foil (alloy 260, Alfa Aesar) was also used to prepare Cu<sub>2</sub>S counter electrodes.<sup>53,54</sup> Brass foil was etched in 37% HCl at 70 °C for 5 min to expose more copper from the copper-zinc matrix, and the yellowish color of the brass changed to a dark red. The etched brass was then rinsed with deionized water and dried at 70 °C. One drop of an aqueous polysulfide solution containing 1 M S and 1 M Na<sub>2</sub>S·9H<sub>2</sub>O was added onto an unmasked part of the etched brass, causing it to suddenly become black, indicating the formation of Cu<sub>2</sub>S. To fabricate SSCs with brass cathodes, a hole was made within 1 mm from the edge of the SnO<sub>2</sub> layer on the photoanode. Photoanodes and cathodes were sealed together in a sandwich configuration using a hot-melt polymer (Surllyn, DuPont). The interelectrode space was filled with an electrolyte by vacuum backfilling. Holes were sealed using a small piece of hot-melt polymer and a microscope coverslip. The electrolyte was composed of 1 M S, 1 M Na<sub>2</sub>S·9H<sub>2</sub>O, and 0.1 M NaOH in deionized water. IPCE spectra were measured with a spectral resolution of ca. 5 nm using a 300 W xenon lamp and a grating monochromator equipped with order sorting filters (Newport/Oriel). The incident photon flux was determined using a calibrated silicon photodiode (Newport/Oriel). Photocurrents were measured using an autoranging current amplifier (Newport/Oriel). Control of the monochromator and recording of photocurrent spectra were performed using a PC running the TRACQ Basic software (Newport). Current-voltage characteristics under simulated AM 1.5 illumination were measured using a Keithley Source Meter and the PVIV software package (Newport). Simulated AM 1.5 illumination was provided by a Newport class A solar simulator, and light intensity was measured using a calibrated Si solar cell. The active area of the cells was defined by a mask to be 0.1199 cm<sup>2</sup>.  $V_{oc}$  intensity characteristics and IS spectra were measured using an Autolab potentiostat/galvanostat and the Nova 1.6 software package. IS experiments were performed with cells under illumination provided by a red LED (center wavelength  $\lambda = 627 \text{ nm}$ ) and biased at the  $V_{oc}$  induced by the illumination. The highest light intensity used was sufficient to produce a  $V_{oc}$  approximately equal to that obtained under AM 1.5 1 Sun illumination, and the incident photon flux was of the order of  $10^{17} \text{ cm}^{-2} \text{ s}^{-1}$ . A 15 mV rms voltage perturbation was used, and the frequency range was  $10^5$  to 0.1 Hz. Different illumination intensities were achieved using neutral density filters mounted in an automated filter wheel system (Newport), which was controlled by the Nova 1.6 software.



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**Supporting Information Available:** Optical characteristics of CdS/CdSe-sensitized TiO<sub>2</sub> electrodes (Figure S1) and equivalent circuit used for analysis of impedance spectra (Figure S2). This material is available free of charge via the Internet at <http://pubs.acs.org>.

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